

7. a) Consider the Gaussian surfaces S_1 and S_2 each of area ΔA in the two slabs (Fig. 11.28).

Let the displacements in the two slabs be \vec{D}_1 and \vec{D}_2 , respectively.

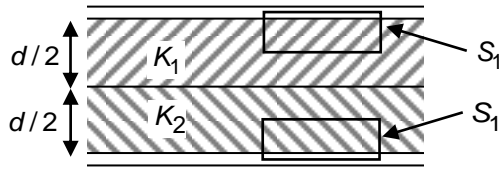


Fig. 11.28: Diagram for answer to TQ 7.

Applying Gauss's law to surface S_1 , we get

$$\int_{S_1} \vec{D}_1 \cdot d\vec{S} = Q_{en}|_{S_1}$$

$$\text{or } D_1 \Delta A = \sigma \Delta A \Rightarrow D_1 = \sigma \quad (\text{i})$$

$$\text{Similarly for surface } S_2, \quad D_2 = \sigma \quad (\text{ii})$$

$$\text{b) Since } \vec{D} = \epsilon_0 K \vec{E}, \text{ we get } E_1 = \frac{D_1}{\epsilon_0 K_1} = \frac{\sigma}{\epsilon_0 K_1} \quad (\text{iii})$$

$$\text{and } E_2 = \frac{D_2}{\epsilon_0 K_2} = \frac{\sigma}{\epsilon_0 K_2} \quad (\text{iv})$$

- c) The potential difference between the plates is given by:

$$V = \int_0^d \vec{E} \cdot d\vec{r} = \int_0^{d/2} \vec{E}_1 \cdot d\vec{r} + \int_{d/2}^d \vec{E}_2 \cdot d\vec{r} = \int_0^{d/2} E_1 dr + \int_{d/2}^d E_2 dr$$

$$\text{or } V = E_1 r \Big|_0^{d/2} + E_2 r \Big|_{d/2}^d = E_1 \frac{d}{2} + E_2 \frac{d}{2}$$

Using Eqs. (iii) and (iv) in this expression, we get

$$V = \frac{\sigma d}{\epsilon_0} \left(\frac{1}{K_1} + \frac{1}{K_2} \right) \quad (\text{v})$$

$$\text{d) From Eq. (v), } C = \frac{Q}{V} = \frac{\sigma A \epsilon_0}{\frac{\sigma d}{\epsilon_0} \left(\frac{1}{K_1} + \frac{1}{K_2} \right)} = \frac{2A \epsilon_0}{d} \frac{K_1 K_2}{(K_1 + K_2)}$$

8. In order to calculate the surface charge densities at $r = b$ and $r = c$, we need to calculate the polarisation for both cases. We do it as follows:
Due to spherical symmetry, the electric fields and the displacements are radial for both cases. Now consider a spherical Gaussian surface of radius r such that $a < r < b$ (Fig. 11.29). Since Q is the charge enclosed by this sphere, from Gauss's law, we have

$$D_1 4\pi r^2 = Q \Rightarrow D_1 = \frac{Q}{4\pi r^2}$$

Now recall from Unit 6 that the electric field

$$\vec{E}_1 = \frac{Q}{4\pi \epsilon_0 r^2} \hat{r} \quad a < r < b \quad (\text{i})$$

In the same way, we can show that for the region $b < r < c$,

$$\vec{E}_2 = \frac{Q}{4\pi \epsilon r^2} \hat{r} \quad b < r < c \quad (\text{ii})$$

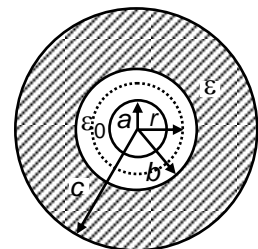


Fig. 11.29: Diagram for answer to TQ 8.

Now, you know that for linear dielectrics,

$$\vec{P} = \vec{D} - \epsilon_0 \vec{E} = (\epsilon - \epsilon_0) \vec{E} = \epsilon_0 (K - 1) \vec{E}$$

Therefore, at $r = b$, $\vec{P}_1 = \epsilon_0 (K - 1) \vec{E}_1$

Since the normal to the dielectric surface at $r = b$ is along $-\hat{r}$, we have

$$\begin{aligned} \sigma_b|_{r=b} &= -\vec{P}_1 \cdot \hat{r}|_{r=b} = -\epsilon_0 (K - 1) \vec{E}_1 \cdot \hat{r}|_{r=b} \\ &= -\epsilon_0 (K - 1) \frac{Q}{4\pi\epsilon_0 r^2} \hat{r} \cdot \hat{r}|_{r=b} \\ &= -\epsilon_0 (K - 1) \frac{Q}{4\pi\epsilon_0 b^2} \end{aligned}$$

or
$$\sigma_b|_{r=b} = -\frac{Q(K-1)}{4\pi b^2}$$

To determine σ_b at $r = c$, we follow the same method as above.

Since the normal to the dielectric surface at $r = c$ (see the margin remark) is along \hat{r} , we have

$$\begin{aligned} \sigma_b|_{r=c} &= \vec{P}_2 \cdot \hat{r}|_{r=c} = \epsilon_0 (K - 1) \vec{E}_2 \cdot \hat{r}|_{r=c} \\ &= \epsilon_0 (K - 1) \frac{Q}{4\pi\epsilon c^2} \\ &= \epsilon_0 (K - 1) \frac{Q}{4\pi\epsilon_0 K c^2} \quad \left(\because K = \frac{\epsilon}{\epsilon_0} \right) \end{aligned}$$

or
$$\sigma_b|_{r=c} = \frac{Q(K-1)}{4\pi K c^2}$$

To determine the potential difference between the outer and inner shells, we begin from its definition

$$V = -\int_c^a \vec{E} \cdot d\vec{r} = \int_a^c \vec{E} \cdot d\vec{r} = \int_a^b \vec{E}_1 \cdot d\vec{r} + \int_b^c \vec{E}_2 \cdot d\vec{r}$$

Substituting \vec{E}_1 and \vec{E}_2 from Eqs. (i) and (ii), we get

$$\begin{aligned} V &= \frac{Q}{4\pi\epsilon_0} \int_a^b \frac{dr}{r^2} + \frac{Q}{4\pi\epsilon_0 K} \int_b^c \frac{dr}{r^2} \quad (\because \epsilon = \epsilon_0 K) \\ &= -\frac{Q}{4\pi\epsilon_0} \left(\frac{1}{r} \right) \Big|_a^b - \frac{Q}{4\pi\epsilon_0 K} \left(\frac{1}{r} \right) \Big|_b^c \\ &= -\frac{Q}{4\pi\epsilon_0} \left(\frac{1}{b} - \frac{1}{a} \right) - \frac{Q}{4\pi\epsilon_0 K} \left(\frac{1}{c} - \frac{1}{b} \right) \end{aligned}$$

Simplifying the expression for V , we get

$$V = \frac{Q}{4\pi\epsilon_0} \left[\frac{c-b}{Kcb} + \frac{b-a}{ab} \right]$$

and
$$C = \frac{Q}{V} = \frac{4\pi\epsilon_0}{\frac{c-b}{Kcb} + \frac{b-a}{ab}}$$

Notice that the unit vector normal to the dielectric's surface points *outward with respect to the dielectric sphere*, which is $+\hat{r}$ at $r = c$ but $-\hat{r}$ at $r = b$.